

## Tutorial Notes # 4. Bohr's Correspondence Principle

For very large values of the quantum number a quantum system will tend to behave like a classical system.

For example, for when  $n$  is large, the particle-in-a-1-d box wave function will have a large number of ~~close~~ closely spaced maxima along the length of the box such that the probability density will tend to become equal everywhere - as for a classical particle. Since  $E = \frac{n^2 h^2}{8ml^2}$ ,

for a macroscopic particle with macroscopic energy ( $m, l, E$  all large),  $n$  will be very large (because  $h$  is very small), so it will behave like a classical particle. No quantisation and ~~spaced~~ uniform distribution within the box. Note

that for if  $\Delta E_n = E_n - E_{n-1}$ , for large  $n$ ,

$$\frac{\Delta E_n}{E_n} \sim \frac{1}{n}$$

Now consider the H atom wave function.

$$\hat{L}^2 \Psi_{n,l,m} = l(l+1) \hbar^2 \Psi_{n,l,m}$$

According to the correspondence principle for very large values of  $l$ ,  $\hbar^2 l(l+1)$  should be identified with the total classical angular momentum of the electron in its orbit.

For  $l = n-1$  and  $m = l$  (highest allowed values for a given  $n$ ), we have

$$\Psi_{n,l,m} \approx \rho^{n-1} e^{-\rho/n} \sin^{n-1}(\theta) e^{i(n-1)\phi}$$

The maximum in the radial part,  $\rho$  (at  $\rho = n(n-1)$ )

as well as the maximum in the  $\theta$  part (at  $\theta = \pi/2$ )

will become more and more sharply defined

as  $n$  increases. Therefore, for very large  $n$

the wave function tend to represent the classical circular Bohr orbit.