relation shows that the emitted radiation for $\Delta M = \pm 1$ is generally elliptically polarized in the (θ, ϕ) plane:

$$\frac{E_\theta^2}{a^2 \cos^2\theta} + \frac{E_\phi^2}{a^2} = 1.$$

It is further seen, that an observer on the northern hemisphere looking toward the source (that is in the direction $-\mathbf{n}$) will recognize right elliptical polarization for $\Delta M = +1$ and left elliptical polarization for $\Delta M = -1$. The results are summarized in Table I.

(*b*) *Intensity and Angular Distribution*

Abbreviating the factor $(c/4\pi)(\omega^2/rc^2)^2$ by μ we find from Eq. (8)

$$I(\Delta M = 0) = \mu 2 \mid \langle J'M \mid p_z \mid JM \rangle \mid^2 \sin^2\theta$$

$$I(\Delta M = \pm 1) = \mu \tfrac{1}{2} \mid \langle J'M \pm 1 \mid p_\pm \mid JM \rangle \mid^2 \times (1 - \cos^2\theta)$$

In observing the Zeeman splitting in a laboratory experiment, θ will equal $\pi/2$ and the intensities are then proportional to the matrix elements only.

[1] W. Heitler, *The Quantum Theory of Radiation* (Oxford U.P., London, 1944).

[2] See e.g., J. L. Powell and B. Crasemann, *Quantum Mechanics* (Addison-Wesley, Reading, MA, 1961).

Addendum to "Modified Hamilton's Principle." John R. Ray [Am. J. Phys. **41,** 1188 (1973)].

Professor R. Weinstock has informed me that the calculation presented in this article also appears in his book, *Calculus of Variations* (McGraw-Hill, New York, 1952), pp. 77–78.

Erratum: "Electron Diffraction at Multiple Slits." Stanley Hirschi [Am. J. Phys. **42,** 4 (1974)].

There are two corrections necessary in our translation of Claus Jönsson's article which originally appeared in *Zeitschrift für Physik* [**161,** 454 (1961)]. In our Fig. 6 the acceleration voltage should be 50 kV. Also, the electron-diffraction photograph and intensity curve of our Fig. 10 should be interchanged with those of Fig. 11.